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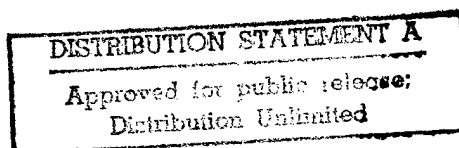
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## Ultrasonic Techniques for the Evaluation of Ceramic Joints

W. A. Simpson, Jr.  
R. W. McClung



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ULTRASONIC TECHNIQUES FOR THE EVALUATION OF CERAMIC JOINTS

W. A. Simpson, Jr., and R. W. McClung

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## ULTRASONIC TECHNIQUES FOR THE EVALUATION OF CERAMIC JOINTS\*

W. A. Simpson, Jr., and R. W. McClung

### ABSTRACT

The increasing use of structural ceramics in high-temperature applications has led to the need for nondestructive evaluation techniques to ensure the integrity of the ceramic materials and the quality of joints consisting of ceramics bonded to ceramics or to metals. We describe the development of ultrasonic techniques for the characterization of ceramic materials and for the detection of flaws in these materials and at ceramic joints. This work has led to the ability to determine which face of a 60- $\mu$ m-thick layer of braze filler material is unbonded, thus providing information about the integrity of the ceramic-filler metal bond. We also describe the development of a rapid technique using Lamb waves to probe the bond between alumina coupons in flexure-strength specimens, whose geometry makes conventional ultrasonic evaluation of the bond difficult.

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### INTRODUCTION

The excellent high-temperature thermal, mechanical, and physical properties of structural ceramics make them likely choices for use in advanced engine designs to allow higher combustion temperatures and therefore higher thermodynamic efficiencies. In addition, the lower weight of such materials relative to that of high-temperature structural alloys should increase an engine's ratio of power to weight. Unfortunately, the low fracture toughness, and hence the small critical flaw size, of structural ceramics precludes the use of standard nondestructive evaluation (NDE) techniques that have been developed over the last 30 years for detection and characterization of critical flaws in metals. For example, flaws of critical size in most structural metals can be detected with ultrasonic waves whose frequencies lie in the range from 1 to 10 MHz. Consequently, most development in ultrasonics has encompassed this frequency range, with little activity occurring above 15 to 20 MHz. In conventional monolithic ceramics, however, the critical flaw size is about

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20  $\mu\text{m}$ , and for such small flaw sizes the use of frequencies of 50 MHz and higher is required. In addition, detection of such small flaws requires the use of focused radiation, and the propagation of such energy through the ceramic surface introduces severe aberration into the beam, thus limiting to a few millimeters the depth at which effective focus can be maintained. Finally, several heat engine designs require that a ceramic be joined to a metal to take advantage of the physical and/or mechanical properties of each. These applications require the development of technology for joining ceramics to other ceramics and to metals and, no less importantly, for inspecting such joints nondestructively to ensure bond integrity. This report describes work performed by the Nondestructive Testing Group of the Metals and Ceramics Division on the development of equipment and techniques for detecting small flaws in ceramic-ceramic and ceramic-metal joints.

The objective of this program is to investigate methods for NDE of ceramic joints leading to recommendations and development of techniques for evaluating important properties and characteristics that affect the serviceability of joints.

## EXPERIMENTAL PROCEDURES AND RESULTS

### CERAMIC MATERIALS

In order to develop techniques for the inspection of ceramic joints it is first necessary to characterize the ceramic materials themselves, since they will often be the determining factor in the choice of test parameters. This is in contrast to the ultrasonic inspection of metals, where the properties other than the wave velocities of the host can frequently be neglected. The engineering constants (Young's modulus, Poisson's ratio, etc.) can easily be determined nondestructively for ceramics by well-known techniques, but this information does not determine the basic inspectability of the sample for a given flaw size. For example, two specimens may have identical engineering constants, but one may be inspectable with 100-MHz ultrasonic energy while the second may not transmit energy above 20 MHz. This behavior results from the fact that

attenuation is highly sensitive to the microstructure of the host; for example, one can glean some information about the average grain size from an attenuation measurement. Therefore, to the determination of the standard engineering constants should be added the attenuation behavior of the particular specimen for ultrasonic energy in a frequency range commensurate with the flaw size of interest.

### Transfer Curve

Although the measurement of attenuation at discrete frequencies is a standard procedure in ultrasonics, the use of such an approach at the high frequencies employed in ceramic inspection would be prohibitively expensive because of the high cost of transducers. A much better approach is to use a single broadband transducer to transmit a range of frequencies and, using a computer, to analyze each frequency component separately. This approach offers the additional advantage of automating the application of corrections to compensate for such nonspecimen losses as beam spread, acoustic impedance mismatch, etc. In addition, the effects of the transducer response and system electronics can be removed from the acquired data to yield a true attenuation versus frequency response for the specimen. This response is termed the transfer curve of the specimen, in analogy to the similar function in optics. In our work, all these corrections have been incorporated into a single computer program that computes the transfer curve from input data consisting of the signals reflected from the front and rear surfaces of a planar sample of the material under test.

When we first began computing transfer curves for ceramics, several anomalies arose that could not be attributed to the samples. In particular, the attenuation in all samples at frequencies above about 30 MHz was much higher than would be expected from optical measurements of the microstructure. This behavior was ultimately traced to the very thin ( $<1\text{-}\mu\text{m}$ ) layer of water used to couple the ultrasonic energy into the ceramic part. For frequencies less than about 15 MHz, the thickness of this coupling layer is such a small fraction of the ultrasonic wavelength that it is negligible. As the frequency increases, however, the coupling layer becomes a larger fraction of a wavelength, and its presence is no longer negligible. When we analyzed this configuration, treating it as a

three-layer rather than a two-layer problem, it became obvious that a correction would have to be added to account for the presence of the coupling layer. Figure 1 shows how the reflection coefficient varies as a function of frequency for a water coupling layer 1  $\mu\text{m}$  thick. At lower frequencies the coefficient approaches the two-layer value, where the effect of the water is negligible. Our computer program was subsequently modified to actually measure the thickness of the coupling layer during a test and provide a correction in the data. The transfer curve so obtained is now highly repeatable and in agreement with destructive analysis.

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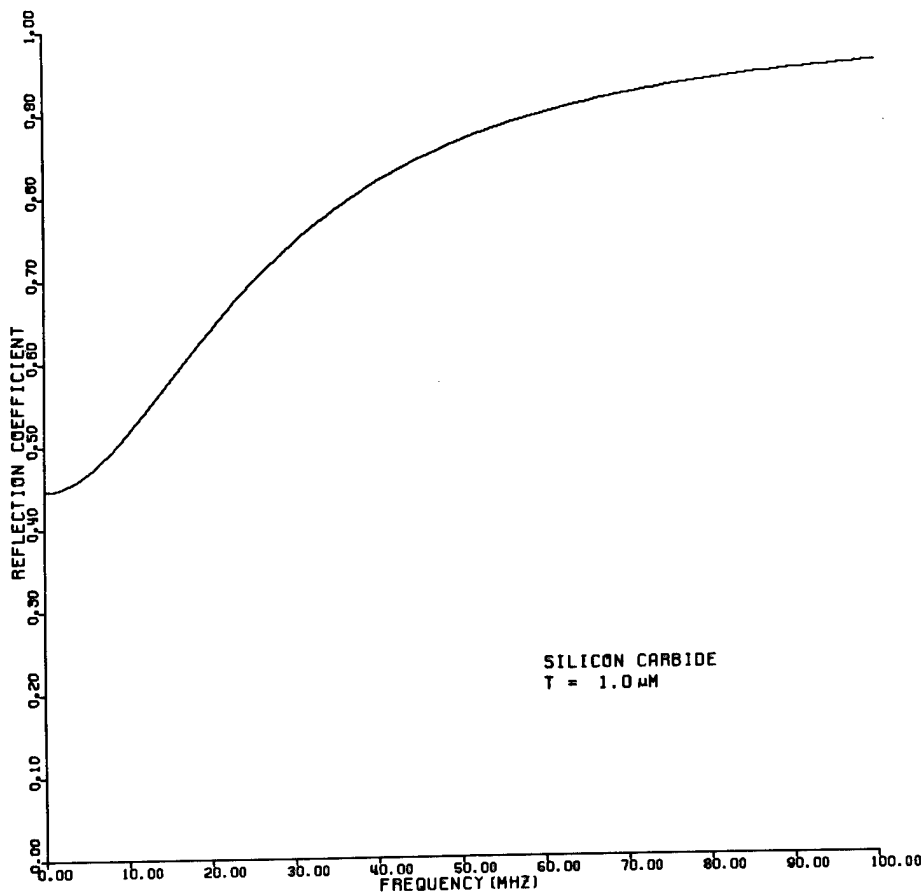


Fig. 1. Reflection coefficient at a water-silicon carbide interface for a 1- $\mu\text{m}$  water layer.

Figures 2 and 3 illustrate the effect of the microstructure on the transfer curve. In Fig. 2, the sample is a piece of tetragonal zirconia polycrystalline (TZP) ceramic with a microstructure of primarily 10- $\mu\text{m}$  tetragonal-phase grains. In Fig. 3, however, the sample is a partially stabilized zirconia (PSZ) with predominantly cubic grains about 100  $\mu\text{m}$  in diameter. It is difficult to propagate through the PSZ an elastic wave whose frequency exceeds about 30 MHz because of severe scattering losses. For such material, the minimum detectable void diameter is probably also about 100  $\mu\text{m}$ .

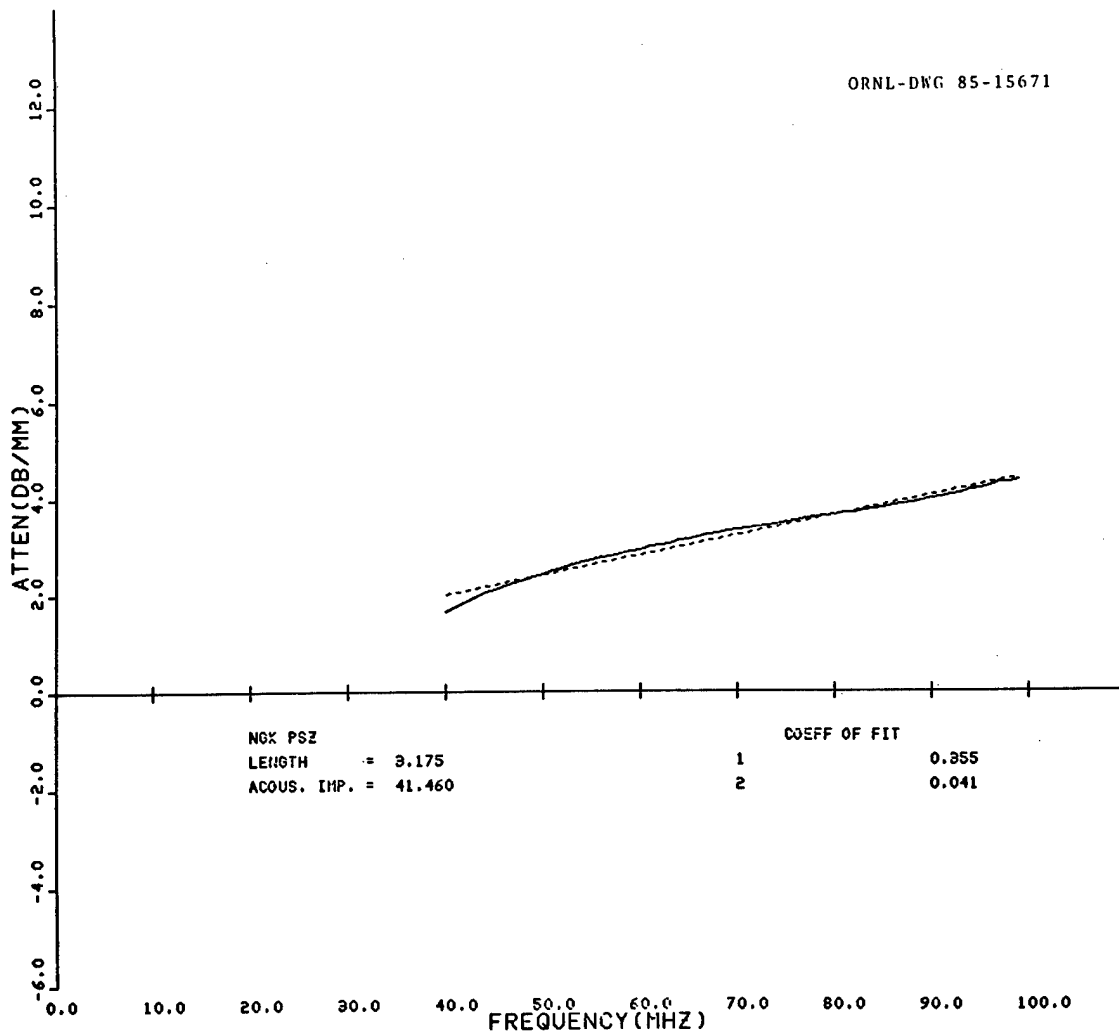


Fig. 2. Transfer curve of tetragonal zirconia polycrystalline ceramic.

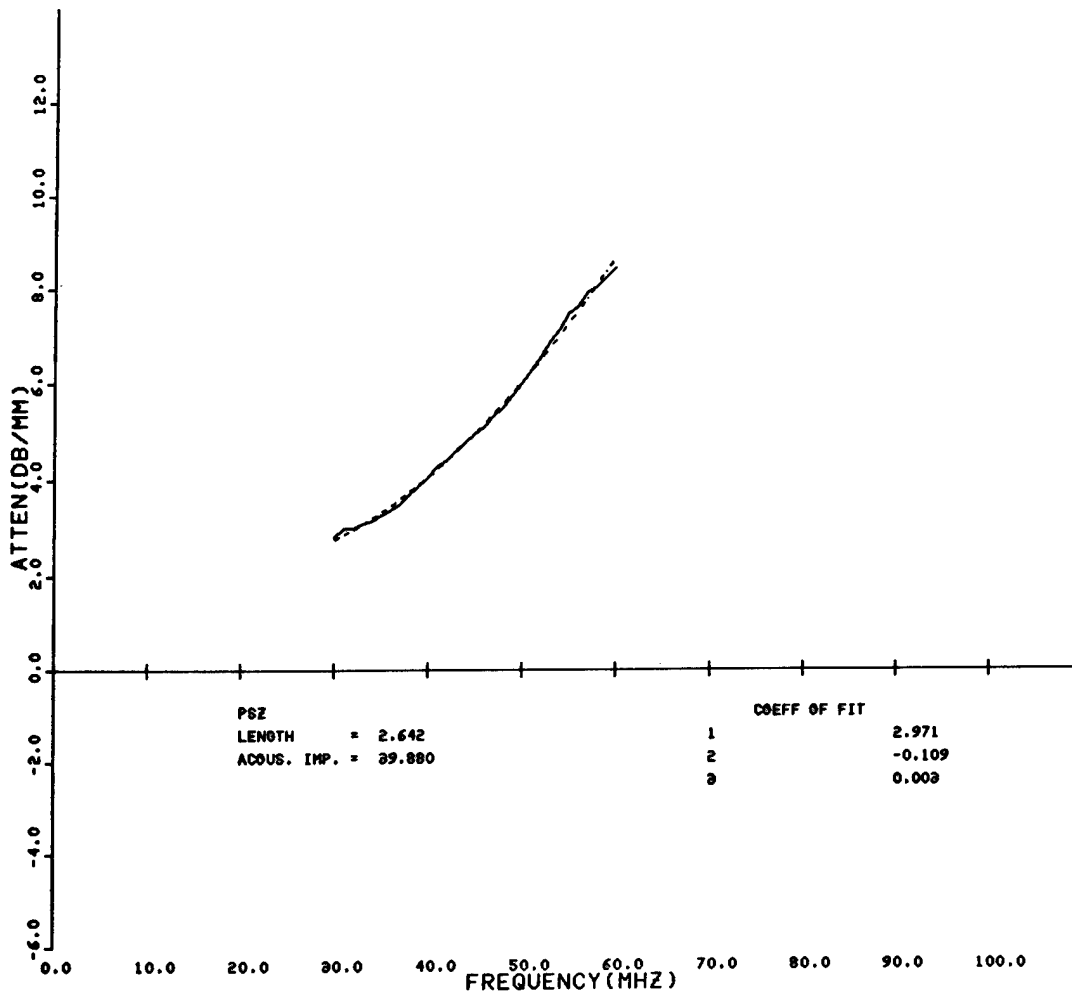


Fig. 3. Transfer curve of partially stabilized zirconia.

#### Detection of Flaws in Ceramics

Also of interest in characterizing the ultrasonic response of a ceramic material is the detection of flaws comparable to or larger than the critical flaw size, particularly if they occur in the vicinity of a joint. From published destructive analyses of flaw initiators in ceramic materials, a common flaw shape is the quasi-spherical void or inclusion. This is fortuitous, since a model exists for the scattering of elastic waves from spherical cavities and inclusions.<sup>1</sup> For a planar crack, we have previously demonstrated a successful model for measuring crack size.<sup>2</sup>

Figure 4 shows the calculated response to elastic waves of a spherical void in silicon nitride. The abscissa is the dimensionless product of the wave number,  $k$  (i.e.,  $2\pi/\lambda$ , where  $\lambda$  is the ultrasonic wavelength in centimeters), of the incident radiation and the void radius,  $a$ . For a given flaw size, the abscissa is thus proportional to frequency. The ordinate is the differential scattering cross section. The response may be divided into three regions: (1) the low-frequency region (Rayleigh scattering) in which the scattering increases as the fourth power of the frequency, (2) the transition region (located at  $ka \approx 1$ ), and (3) the high-frequency region in which the scattering cross section is oscillatory with an approximately constant average value. The latter behavior can be traced to interference between two waves scattered by the sphere. The first is the direct wave reflected from the near point of the sphere, and the second is the so-called "creeping" wave, which, after tangential incidence, propagates around the sphere and is reemitted after

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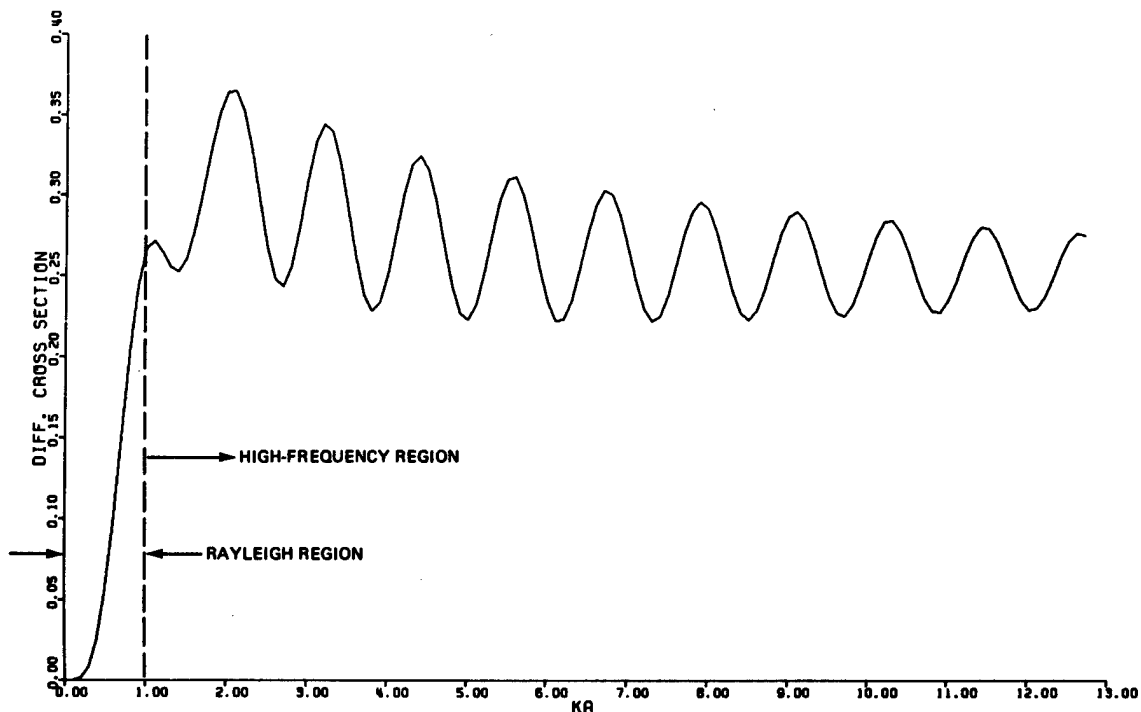


Fig. 4. Differential cross section for scattering from a spherical void.

traversing half the circumference. The periodicity of the oscillation is thus related to the size of the void. For natural flaws, however, even those whose shape is very nearly spherical, the surface roughness may rapidly attenuate the latter wave, and the high-frequency scattering depicted in Fig. 4 may be replaced by a near-constant value. In this case, the size of the scattering center may be approximated by noting that the transition region occurs at  $ka \approx 1$ ; thus if the frequency of the turning point, that is, the transition from Rayleigh to high-frequency scattering, is known, the radius of the sphere may be estimated.

Figure 5 shows experimental data obtained on a natural flaw in TZP ceramic. Here the turning point occurs at about 89 MHz. Based on the measured velocity, the flaw diameter was estimated to be approximately 25  $\mu\text{m}$ . That the flaw size is of this order of magnitude is also supported by the fact that it could not be detected reliably with a scan increment of 50  $\mu\text{m}$  but could with an increment of 25  $\mu\text{m}$ .

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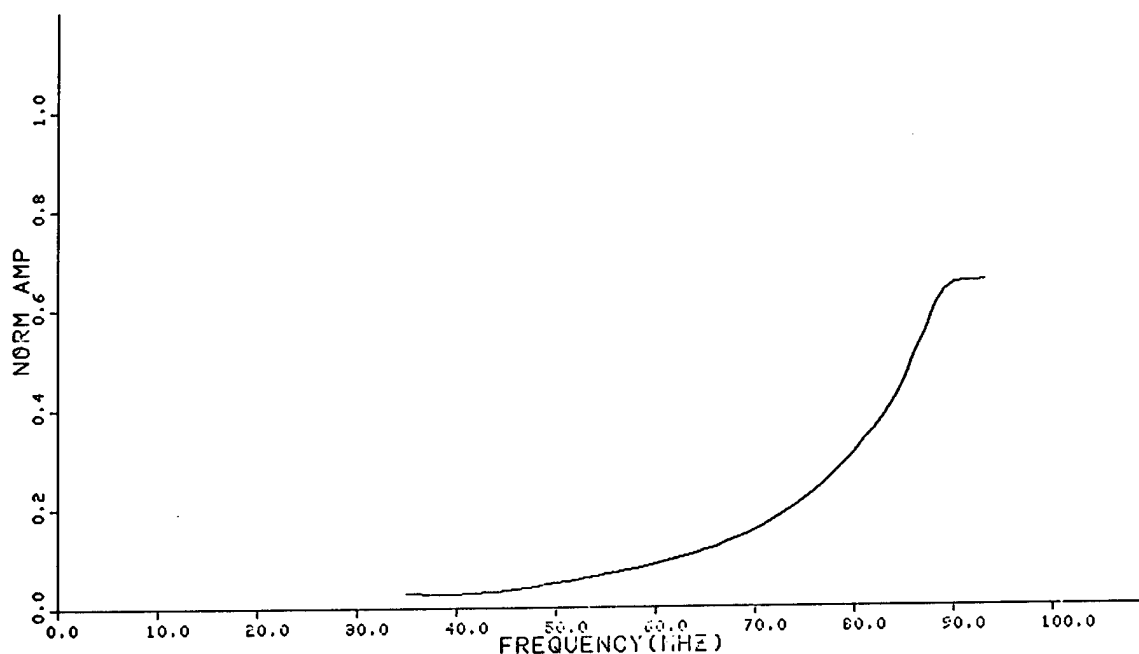


Fig. 5. Experimental data for scattering from a natural flaw in tetragonal zirconia polycrystalline ceramic.

## PLANAR JOINTS

Ceramic-Cap Piston Specimens

Two large ceramic-to-metal brazements that mock up the attachment of a ceramic cap to a diesel engine piston were made available for nondestructive testing studies. These specimens are 111 mm (4 3/8 in.) in diameter and consist of a 6.4-mm-thick (0.25-in.) PSZ cap brazed to a nodular cast iron (NCI) disk via a titanium transition piece and a commercial Ag-Cu-Sn brazing filler metal.<sup>3</sup> The surface of the ceramic had been vapor coated with titanium to promote wetting by the filler metal. Figure 6 shows the geometry of the joint and Fig. 7 one of the specimens. The thickness of each braze layer was about 60  $\mu$ m and that of the titanium transition piece about 0.6 mm.

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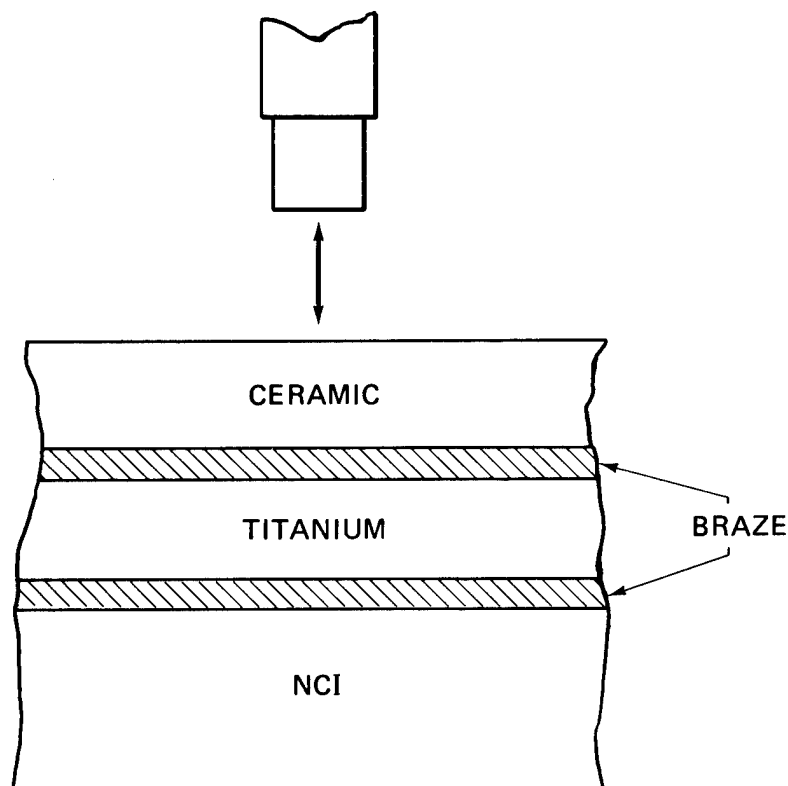


Fig. 6. Diagram of a transition joint for bonding a ceramic to nodular cast iron (NCI) to simulate a diesel-engine piston with a ceramic cap.

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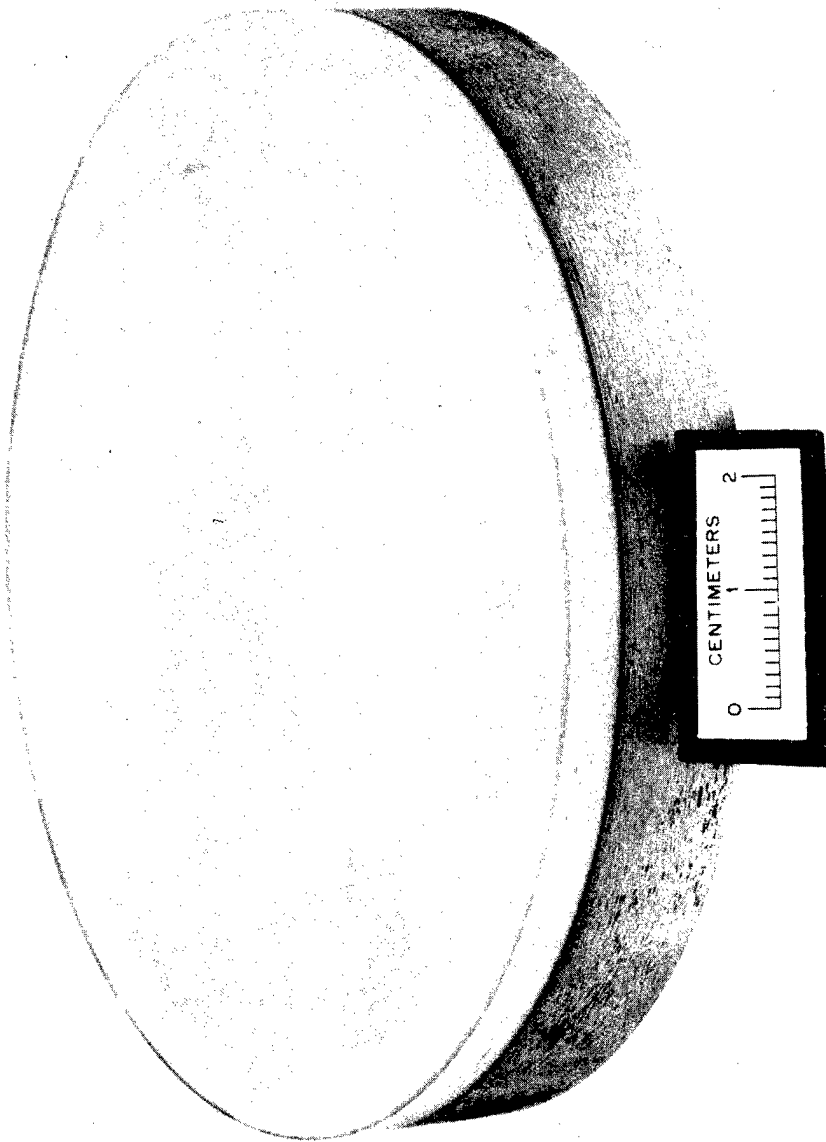


Fig. 7. Ceramic-cap diesel-engine piston specimen (one of two).

Initial ultrasonic examination of each specimen indicated a relatively high attenuation for elastic waves in the ceramic cap, restricting the maximum usable frequency to about 25 MHz. From our previous studies of the attenuation characteristics of PSZ material, this behavior was consistent with the scattering losses produced by the relatively large grain size ( $\approx 100 \mu\text{m}$ ) of this ceramic.

Both specimens were scanned in our high-frequency (100-MHz) ultrasonic system, which permits variations in the transmission of the ceramic-metal bond to be displayed as a gray-scale image. A flat (unfocused), broadband transducer was used, and both samples were found to contain nonbonded regions near the center. Figure 8 shows the result for one of the brazements. Here lighter areas indicate relatively better transmission through (less reflection from) the bond, and darker areas indicate relatively poorer transition (greater reflection). The dark area near the center is completely unbonded. (In the original data this area was uniformly dark; it did not reproduce well photographically, however.) The dark ring near the periphery of the sample is an edge effect caused by the unfocused transducer.

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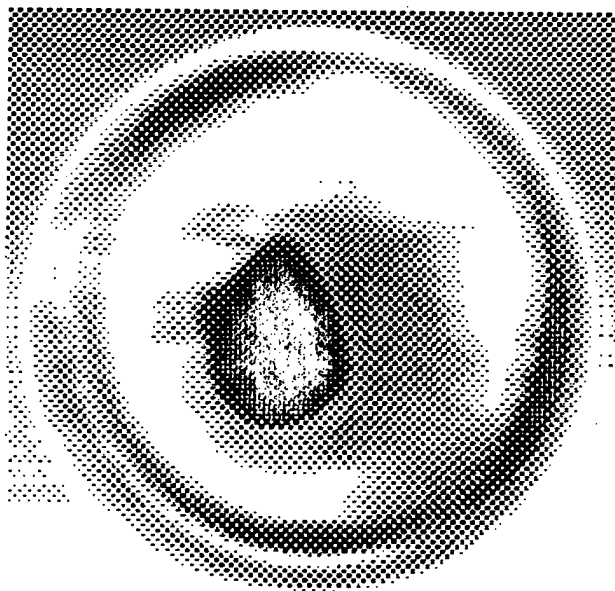


Fig. 8. Gray-scale presentation of ultrasonic transmission through the bond in piston specimen 1.

The samples were scanned a second time using a focused transducer. Focusing minimizes edge effects and more sharply delineates the region of unbond. Figure 9 shows the raw gray-scale data for the second sample, which is seen to be unbonded in the center and at several locations around the periphery.

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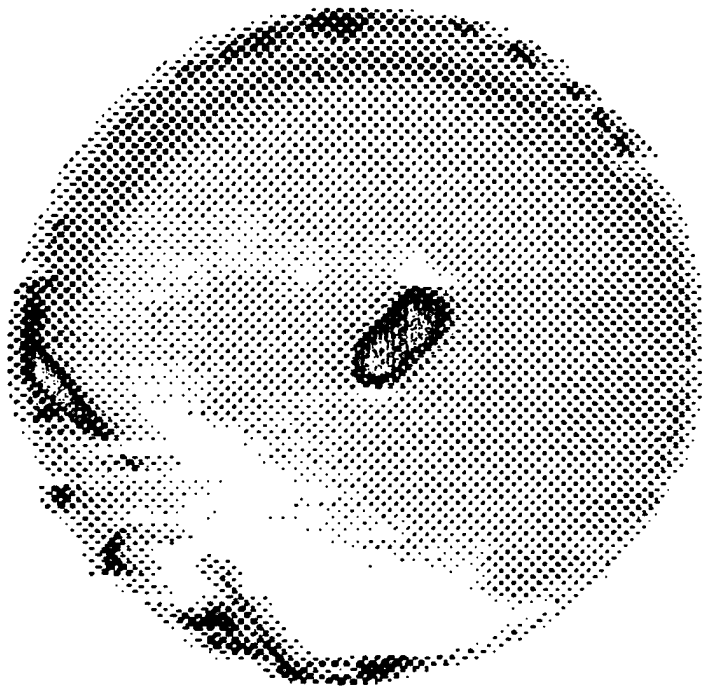


Fig. 9. Gray-scale presentation of ultrasonic transmission through the bond in piston specimen 2.

The raw pixel data from these gray-scale images were next computer processed to permit expansion of a selected range of gray-scale values. Such processing allows minor ( $<3$ -dB) bond variations to be ignored and fluctuations within the unbonded areas to be enhanced. Figures 10 and 11 show the processed images for the two specimens. As expected, the areas delineated show virtually no variations, which is indicative of complete unbonding. The shape, size, and location of the affected areas are also precisely determined by this processing.

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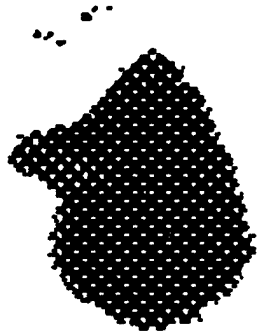


Fig. 10. Enhanced gray-scale presentation of ultrasonic transmission through the bond in piston specimen 1.

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Fig. 11. Enhanced gray-scale presentation of ultrasonic transmission through the bond in piston specimen 2.

Because there are several interfaces in the PSZ-NCI joint (Fig. 6), one would like to know at which of these interfaces the unbond occurs, since this information could provide insight into the origin of the problem. In addition, while the precise location of the unbond may be immaterial from an accept-reject point of view, such information is important in the development of ceramic brazing technology. For example, if the unbond occurs at the ceramic-braze interface, it could indicate a failure of the braze filler metal to wet the ceramic surface, and wettability is a major consideration of the brazing development program. Fortunately, the location of the unbond can be determined nondestructively from the ultrasonic scattering data. Figure 12 shows the rf waveform of signals scattered by the transition joint in three areas of the specimen whose data were shown in Figs. 9 and 11. In Fig. 12, the upper trace (trace A) is that obtained from a well-bonded region. The two signals are from the faces of the 600- $\mu$ m-thick titanium transition piece, with the PSZ to the left and NCI to the right. The ultrasonic wave is incident from the PSZ. In trace B, the second signal is much larger, and the third signal, resulting from reverberation in the titanium, indicates that the unbond occurs on the NCI side of the titanium. In trace C, the first signal is much larger, and the absence of the signal from the opposite face of the titanium indicates that the sample is unbonded on the PSZ side of the titanium.

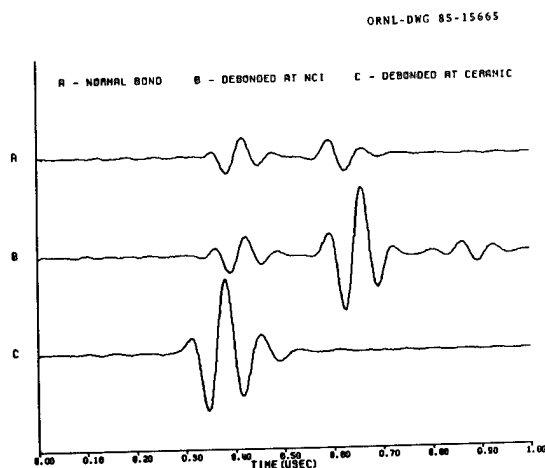


Fig. 12. Radio-frequency waveform of the ultrasonic signal scattered by the interface of piston specimen 2 for various bond conditions.

If the ceramic material will support elastic waves of a sufficiently high frequency, the signals from opposite faces of the nominal 60- $\mu\text{m}$ -thick braze layer can be resolved (a time resolution of only 30 ns), and it is then possible to tell directly whether the unbond occurs at the ceramic-braze or braze-titanium interface. For example, Fig. 13 shows the waveform obtained from the joint region of a sample containing TZP (wide bandwidth) ceramic. The first two signals originate at the faces of the nominal 60- $\mu\text{m}$ -thick braze layer between the TZP and the titanium transition piece. The second two are from the similar braze layer between the titanium and the NCI substrate. Note that the signals from the braze layers are resolved, which permits the condition of any of the four interfaces to be monitored. The restricted bandwidth of the PSZ is not sufficient to resolve these signals (frequencies above about 30 MHz are not transmitted). However, we have recently developed a technique to enhance the resolution of closely spaced ultrasonic signals.<sup>4</sup> Use of this technique will generally produce delta functions in the output data at the location of each individual wave center in the input data, even when the waves are too closely spaced to resolve in the time domain. When this technique was used to process the signals from the data of Fig. 12, the signals from either face of the braze layers could be resolved, as shown in Fig. 14. Since all four signals are present, the region is one in which the sample is well bonded. For the central unbond region of the simulator shown in Fig. 11, however, the processed data show only a single

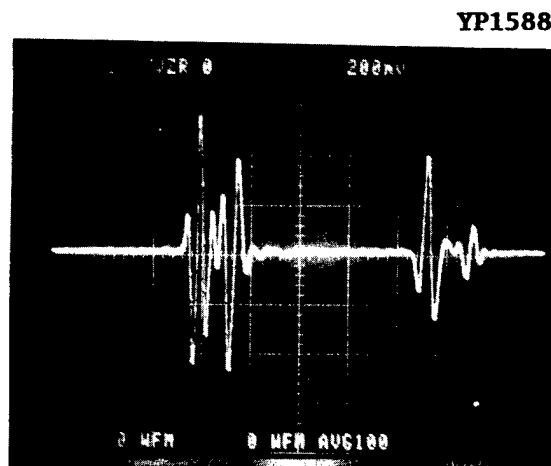


Fig. 13. Radio-frequency waveform of the ultrasonic signal scattered by the interface in tetragonal zirconia polycrystalline ceramic.

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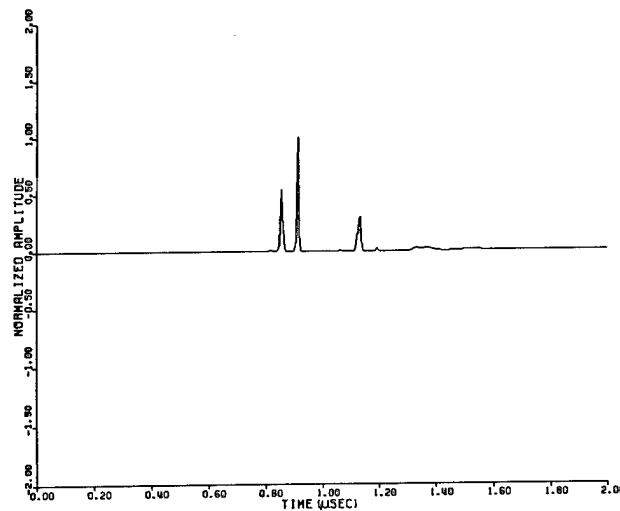


Fig 14. Processed data showing recovery of four interface signals in the bonded region of tetragonal zirconia polycrystalline ceramic.

delta function at the location of the PSZ-braze interface, indicating that the problem is failure of the braze to adhere to the ceramic.

These results indicate that for both samples, the central unbond occurs at the PSZ interface. For the specimen of Figs. 9 and 11, some of the edge indications are unbonded at the PSZ interface and some at the NCI interface.

Following NDE, the Materials Joining Group sectioned the sample of Figs. 9 and 11 through the unbonded region delineated by the central indication. When the region was cut out, the ceramic cap actually fell off. Subsequent metallographic examination of the ceramic-braze interface revealed the presence of extreme porosity, possibly due to trapped gas. The nondestructive results were thus validated.

#### Shear Specimens

Following examination of the piston-cap specimens, some smaller ceramic-metal brazements were made available. These samples were approximately 9-mm-square pieces of 3.5-mm-thick zirconia ceramic brazed to an NCI substrate with a 0.6-mm-thick titanium transition piece. In this case, however, the ceramic was a fine-grained TZP material that could support elastic waves at frequencies in excess of 100 MHz. In such

samples, effective focus could be maintained at depths up to 5 mm for these frequencies: accordingly, a 100-MHz focused transducer with a focal length of 25.4 mm in water was obtained to exploit this condition. At this frequency and with the transducer focused at the PSZ-braze interface, the signals from either side of the 60- $\mu\text{m}$  braze layer could be resolved, even though these signals are separated by only about 30 ns. Therefore, we were able to gate selectively the signal from the TZP-braze interface for analysis. This signal is an indicator of the quality of the bond that exists between the braze filler metal and the TZP. Figure 15 shows two of the samples and Fig. 16 the results of scanning the bond region. In these gray-scale presentations, lighter areas indicate relatively better bonding while darker regions depict relatively poorer bonding. The sample on the right has numerous small regions where there is lack of bonding. These regions average perhaps 100  $\mu\text{m}$  in diameter and are probably caused by bubbles of trapped gas in the braze material. The presence of these bubbles has been previously demonstrated by destructive analysis and inferred from examination of the fracture surface, but, until now, we have not been able to detect their presence nondestructively.

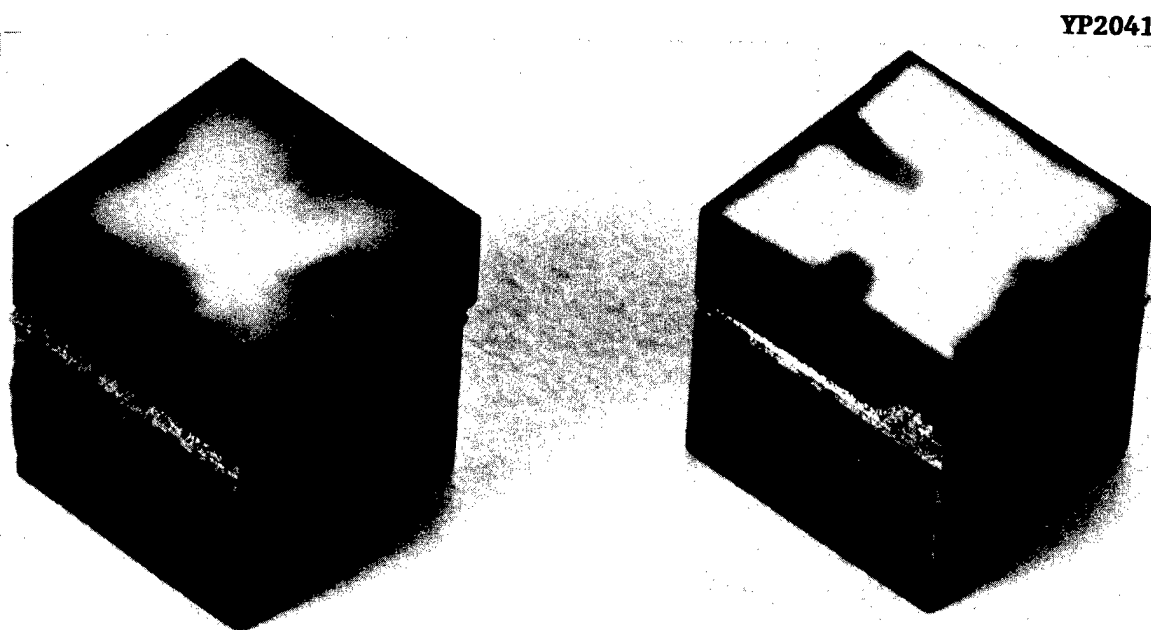


Fig. 15. Tetragonal zirconia polycrystalline ceramic joint specimens.

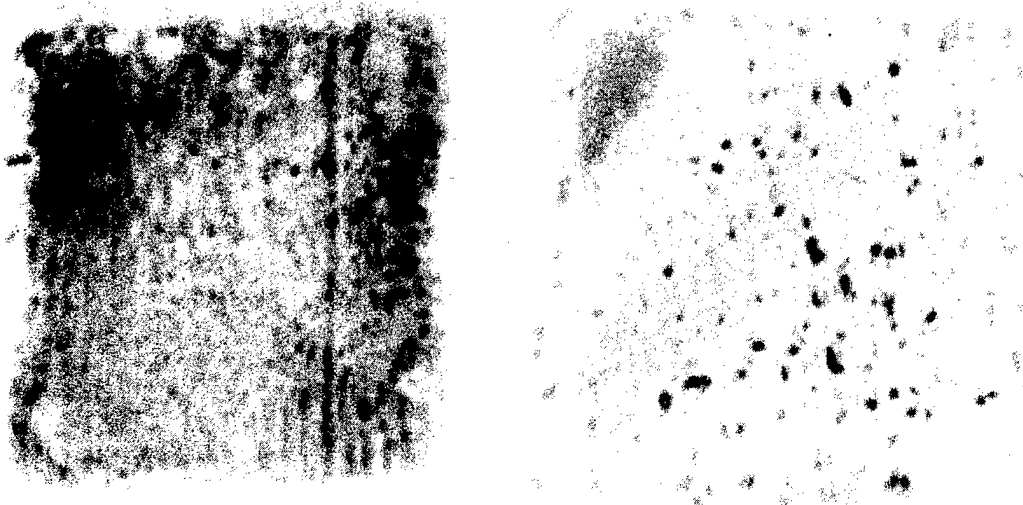


Fig. 16. Gray-scale presentation showing the detection of pores in the braze layers of the ceramic joint specimens shown in Fig. 15.

A diagonal band, running from lower left to upper right of the sample on the right, can be seen in the data of Fig. 16. The cause of this darkening is not known, but possibilities include the presence of microporosity or residual stress in the braze layer. The dark area in the upper left of the sample on the right is caused by severe thinning of the braze material.

The generally darker nature of the sample on the left in Fig. 16 is possibly attributable to use of a different grade of ceramic. Small variations in the acoustic properties alter the reflection coefficient at the TZP-braze interface and vary the average brightness of the reproduced data.

The sample on the left also exhibits a periodic variation in the interface that may be caused by machining marks on the TZP. No such marks are detectable on the visible surface, but they sometimes occur on one or more surfaces of the blanks.

The above results indicate that a great deal of information about the nature of the PSZ-braze bond can be gleaned from the ultrasonic scattering if the signals from the braze layer can be resolved. However, if they cannot, as is the case at typical ultrasonic frequencies (1-10 MHz), the signal generated by subtle variations at the filler metal-PSZ interface will be swamped by the signals from the filler metal-titanium interface.

#### Manufactured Flaws

We have considerable interest in determining the minimum detectable area of unbond in a ceramic joint. From the backscattering spectrum obtained from discrete flaws in ceramics, we earlier showed (see Fig. 5) the ability to infer a minimum measurable flaw size of about 25  $\mu\text{m}$  (the minimum detectable flaw size will be still smaller) with our present system. It is difficult to fabricate discontinuities in this size range. However, we obtained a standard consisting of a bonded couple between a P-leg and an N-leg of a silicon-germanium thermoelectric sample. The bond contains three manufactured flaws having diameters of 250, 650, and 125  $\mu\text{m}$ , and the acoustic properties (i.e., velocity and density) of the silicon-germanium are similar to those of typical ceramics. Figure 17

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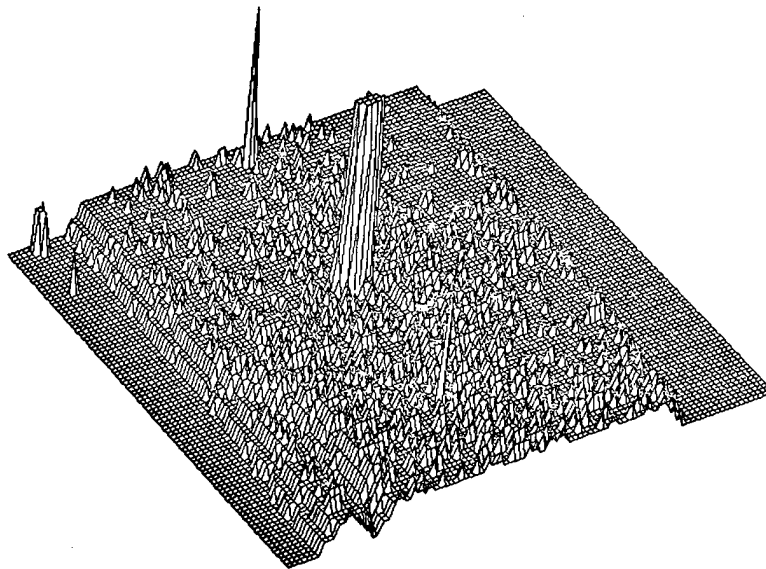


Fig. 17. Three-dimensional presentation of the ultrasonic detection of, left to right, 250-, 635-, and 125- $\mu\text{m}$  flaws at the interface between two silicon-germanium thermoelectric samples.

is a pseudo three-dimensional view of the ultrasonic scattering from the bond line showing the detection of all three flaws. In the lower left of the figure is a natural crack that extended through the substrate and terminated at the interface. Figure 18 is a view of the data from a lower angle, making the smallest flaw more visible. Figure 19 is an expansion of the region around the 125- $\mu\text{m}$  flaw. Note that the flaw signal is much larger than the background from the interface; thus, we should be able to see flaws considerably smaller than 125  $\mu\text{m}$  reliably.

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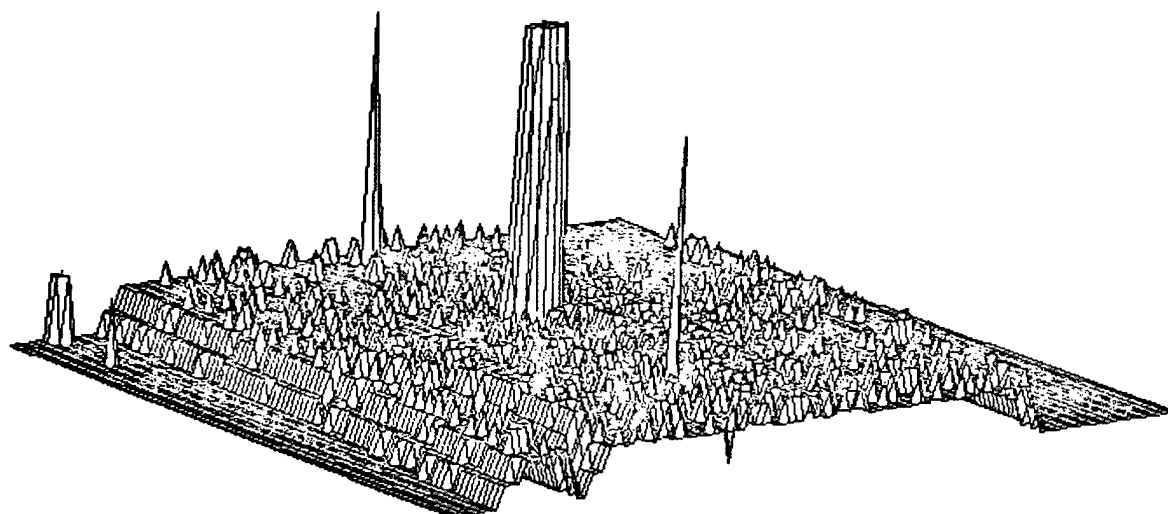


Fig. 18. Three-dimensional presentation (from a lower angle than in Fig. 17) of the ultrasonic detection of, left to right, 250-, 635-, and 125- $\mu\text{m}$  flaws at the interface between the two silicon-germanium thermoelectric samples.

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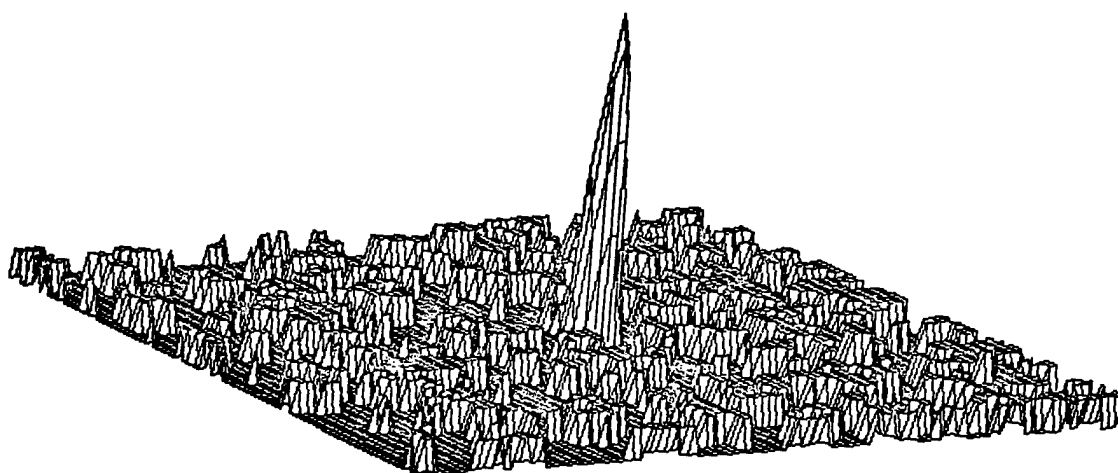


Fig. 19. Three-dimensional presentation of the region around the 125- $\mu\text{m}$  flaw at the interface between the two silicon-germanium thermoelectric samples.

## BUTT JOINTS

### Lamb Wave Studies

For the samples described earlier in this report, the specimen geometry was such that bulk-wave probing of the joint region (albeit at frequencies far above those used in conventional ultrasonic testing) could be performed at normal incidence. We demonstrated earlier the ability to resolve signals from either side of the 60- $\mu\text{m}$ -thick layer of brazing filler metal when the ultrasonic waves were incident normally on the joint. This capability then allowed us to determine whether a lack of bond occurred at the ceramic or metal interface of the joint and to measure the variation in filler metal thickness across the specimen.

For some specimen configurations, normal-incidence testing of the braze region using bulk waves is not possible. For example, a butt-braze joint between thin plates requires angle beams if bulk waves are used to interrogate the interface. While this configuration may or may not be important in the final application of ceramic joining to heat engines, it is of considerable importance to the development of joining technology in that many test specimens use such joint geometries. For example, specimens to measure flexural strength of ceramic-ceramic or ceramic-metal brazements are of this type (see Fig. 20), and inspection techniques must be developed so that correlations between NDE indications and mechanical properties can be determined. Only in this way can one differentiate between significant and ignorable indications.

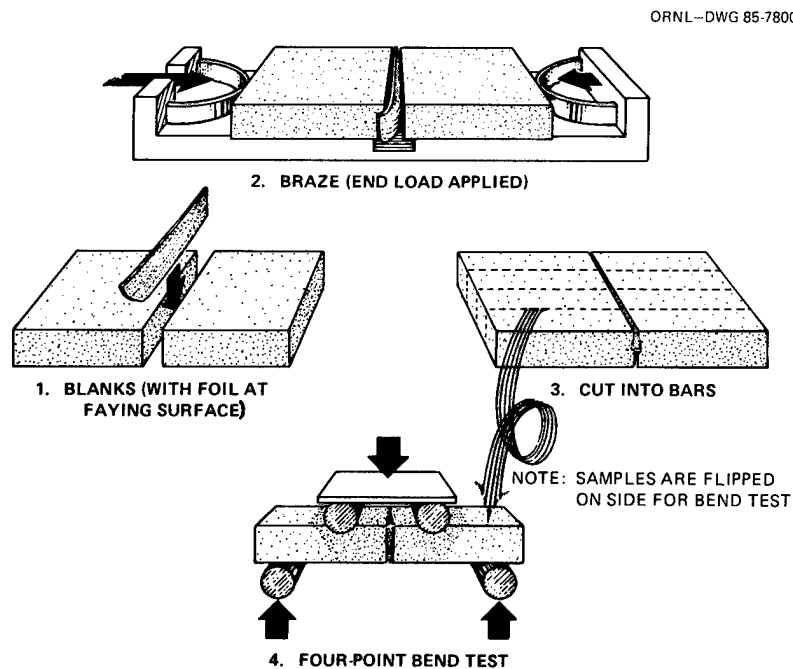


Fig. 20. Fabrication of ceramic flexure-strength specimens with butt-braze joints.

We have examined several butt-braze joints in  $25.4 \times 14.2 \times 2.9$  mm alumina coupons using high-frequency angle-beam techniques. While this approach would appear to have considerable potential for evaluating bond quality, the results thus far have been more difficult to interpret than the results from normal-incidence tests, and they suffer from the fact that the focal point of the transducer cannot be maintained exactly on the bond across the full height of the joint. This test is also relatively slow, since the bond must be scanned using a small linear increment.

A second approach for evaluating bond quality in butt-joint specimens relies on measurement of the transmission coefficient of so-called plate or Lamb waves propagated through the bond. In this technique, a Lamb wave, which is the elastomechanical analog of guided electromagnetic waves, is excited on one side of the joint and the amplitude of the wave on the second side measured after propagation through the bond. This approach has the advantage of requiring scanning only along a single axis, making it much faster than conventional testing. The frequencies involved are also quite low, typically 1 to 5 MHz. The quantity determined by this measurement is the relative area of the bonded region, which has been related to the prediction of shear strength for spot welds in metals.<sup>5</sup>

As is the case for electromagnetic guided waves, Lamb waves are highly dispersive. Since the propagation of such a wave along a plate produces (microscopic) flexure of the plate, Lamb waves can produce either symmetric or antisymmetric modes according to the symmetry of the flexure with respect to the center line of the plate. Figures 21 and 22 show the calculated dispersion curves for the first four symmetric and antisymmetric modes in 2.9-mm-thick alumina. Since the ordinate gives the Lamb wave phase velocity (normalized by the shear wave velocity), it is also related to the angle of incidence necessary to generate the given Lamb mode. Thus a horizontal line will intersect the various possible modes, whose abscissas give the frequencies necessary to establish the modes at the indicated angle of incidence of the exciting energy.

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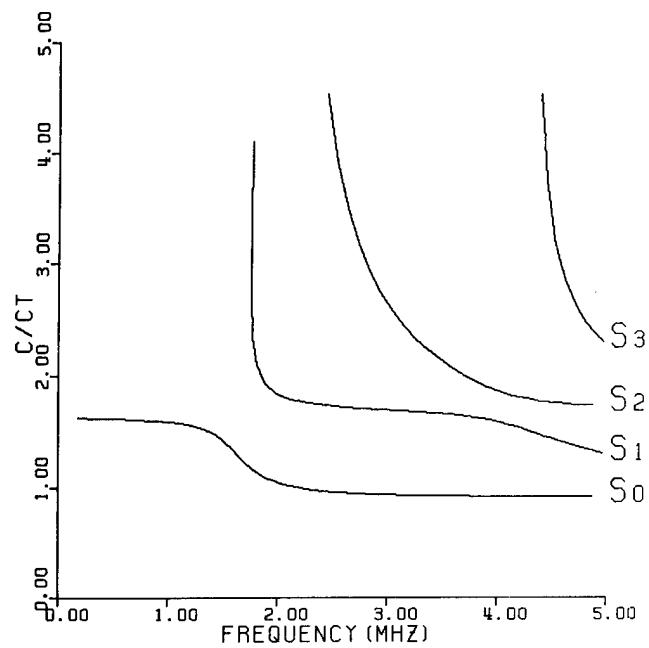


Fig. 21. Dispersion curves for symmetric Lamb wave modes in alumina.

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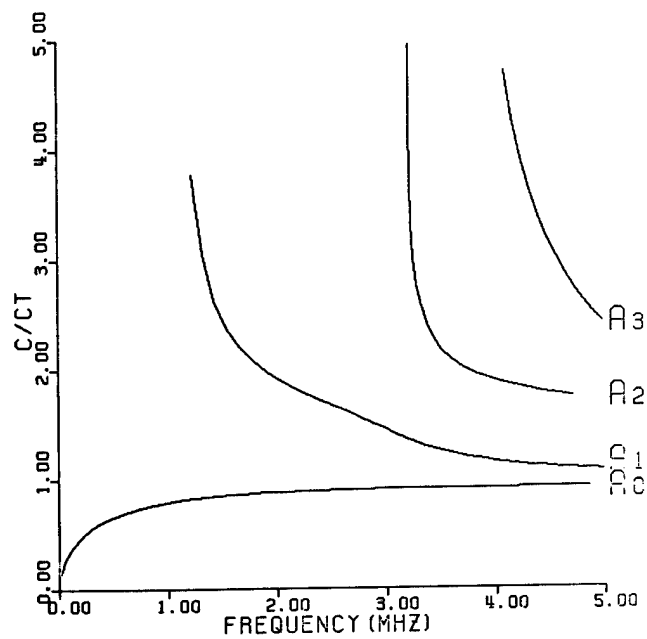


Fig. 22. Dispersion curves for antisymmetric Lamb wave modes in alumina.

Figure 23 is an experimental result showing the excitation of five Lamb wave modes in a 2.9-mm-thick alumina coupon. The normalized velocity of each wave is 1.25, and the mode frequencies are in good agreement with the values shown in Figs. 21 and 22. Note that this normalized velocity does not intersect the dispersion curve of the lowest-order antisymmetric mode; thus, this wave is not present in the spectrum. The relative amplitudes of the various waves are determined primarily by the response of the exciting transducer, which, in these studies, is a broadband 5-MHz unit.

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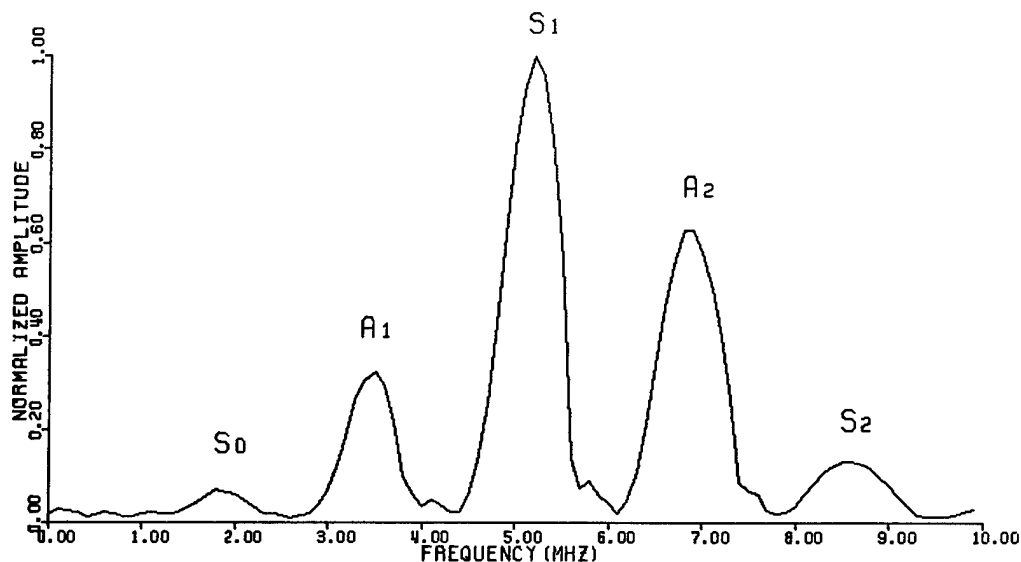


Fig. 23. Spectrum of Lamb wave modes in an alumina coupon.

A pure-mode response can be obtained by using single-frequency excitation of an appropriate transducer. Alternatively, one can use a relatively narrow-band transducer and select a normalized velocity (incident angle) such that only one mode exists within the passband of the transducer. Figure 24 shows the excitation of a single mode, the lowest-order symmetric mode, in an alumina coupon by a 2.25-MHz transducer using pulse excitation.

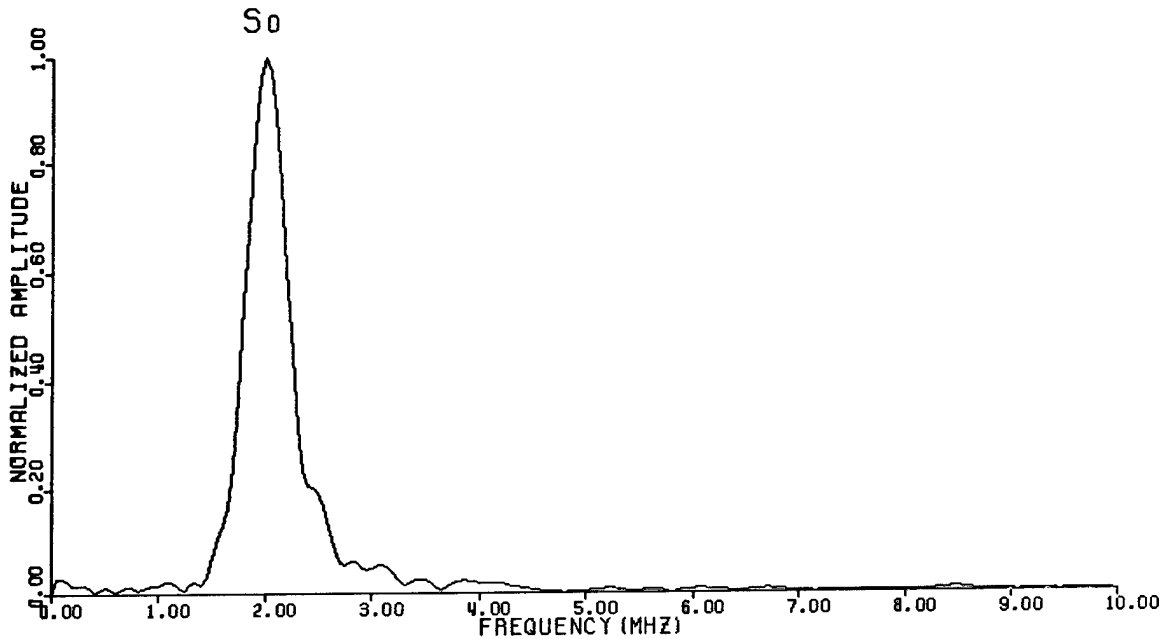


Fig. 24. Excitation of a single Lamb mode in an alumina coupon.

In order to test the effectiveness of using Lamb waves to interrogate the butt-braze joint in alumina flexure-strength specimens, the specially fabricated sample shown in Fig. 25 was made. A piece of tantalum foil at each end of the joint provides the proper separation of the ceramic halves during melting of the braze. A third piece of foil was placed in the center of this particular specimen to prevent bonding in that region and to simulate a nonbond of known dimensions. The brazing filler metal was one of several under development at Oak Ridge National Laboratory that will wet oxide ceramics directly with no pretreatment of the ceramic surface.

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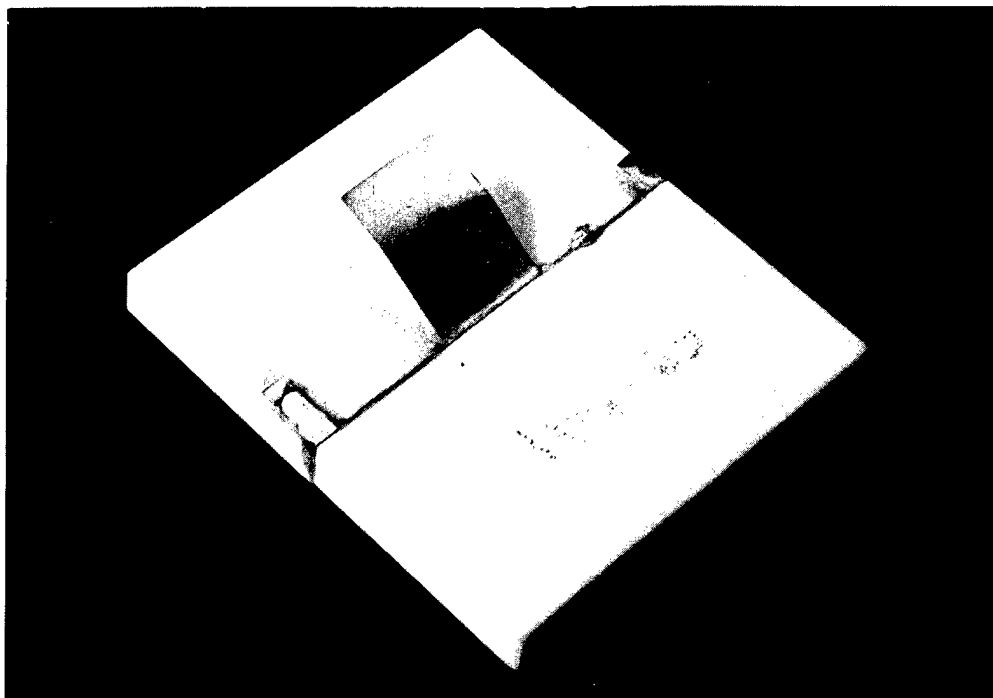


Fig. 25. Flexure-strength specimen containing a simulated nonbond.

The standard was first examined by conventional angle-beam through-transmission techniques using 25-MHz transducers. The joint was scanned in an x-y pattern using an index of 100  $\mu\text{m}$ . As expected, the central nonbond was easily detected using both flat and focused transducers. For the latter, the sensitivity varied from top to bottom of the interface because of variations in the beam profile with depth. In both cases, however, numerous indications were generated at the top and bottom of the interface by a slight vertical misalignment of the ceramic coupons. This particular approach is sensitive to misalignment and is a disadvantage of angle-beam testing of butt joints.

We next examined the standard using Lamb waves. Two transducers were used with the transmitter and receiver located on opposite sides of the interface. Figure 26 shows the transducer configuration. The transmitter was driven with a tone burst to ensure excitation of a single Lamb mode.

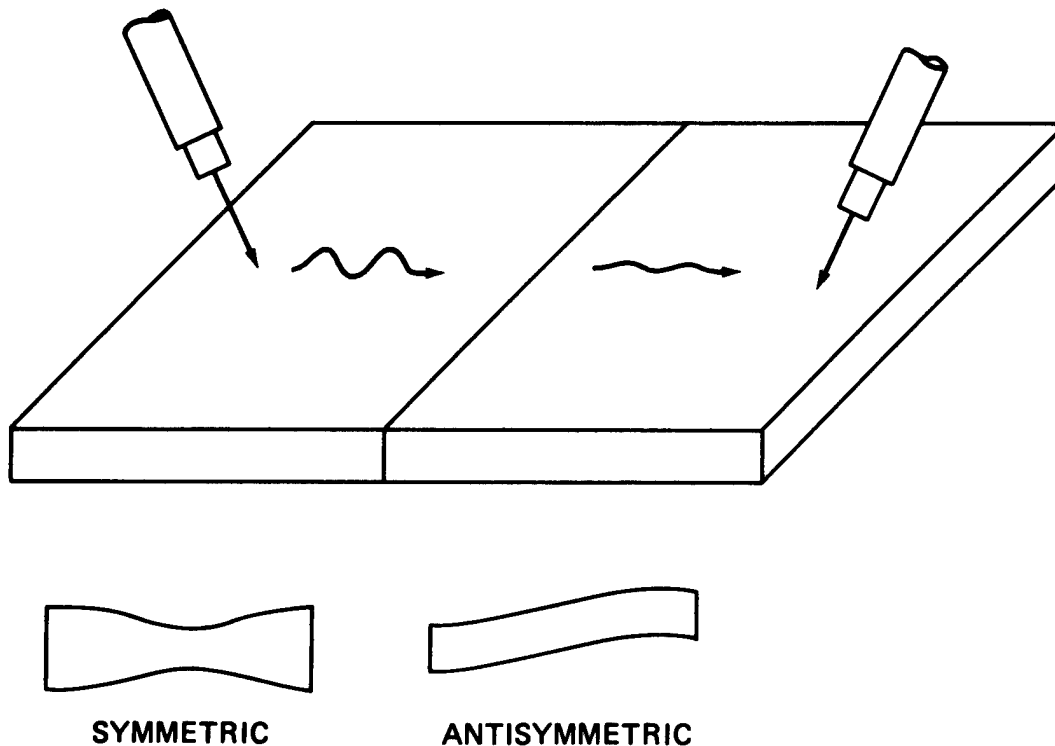


Fig. 26. Transducer configuration for excitation and detection of Lamb waves in the butt-brazed alumina coupon illustrated in Fig. 25.

As expected, virtually no change in transmission amplitude was detected when the transducers were translated perpendicularly to the joint. Nevertheless, an x-y scan (parallel and perpendicular to the interface) was made. The results are shown in Fig. 27, where the height of the surface represents the Lamb wave transmission amplitude. The central nonbond corresponds to the large dip in the surface. Obviously, a single scan parallel to the interface would have sufficed, so this test can be performed very rapidly.

A second butt-braze specimen was also available; it was a standard flexure bar consisting of titanium-coated PSZ coupons brazed with a commercial Ag-Cu-Sn filler metal. Only the two end pieces of tantalum foil were present in this sample. It was first tested by conventional angle-beam techniques with no indications detected other than the usual ones attributable to vertical misalignment of the ceramic plates. In particular, no regions of unbonding were found.

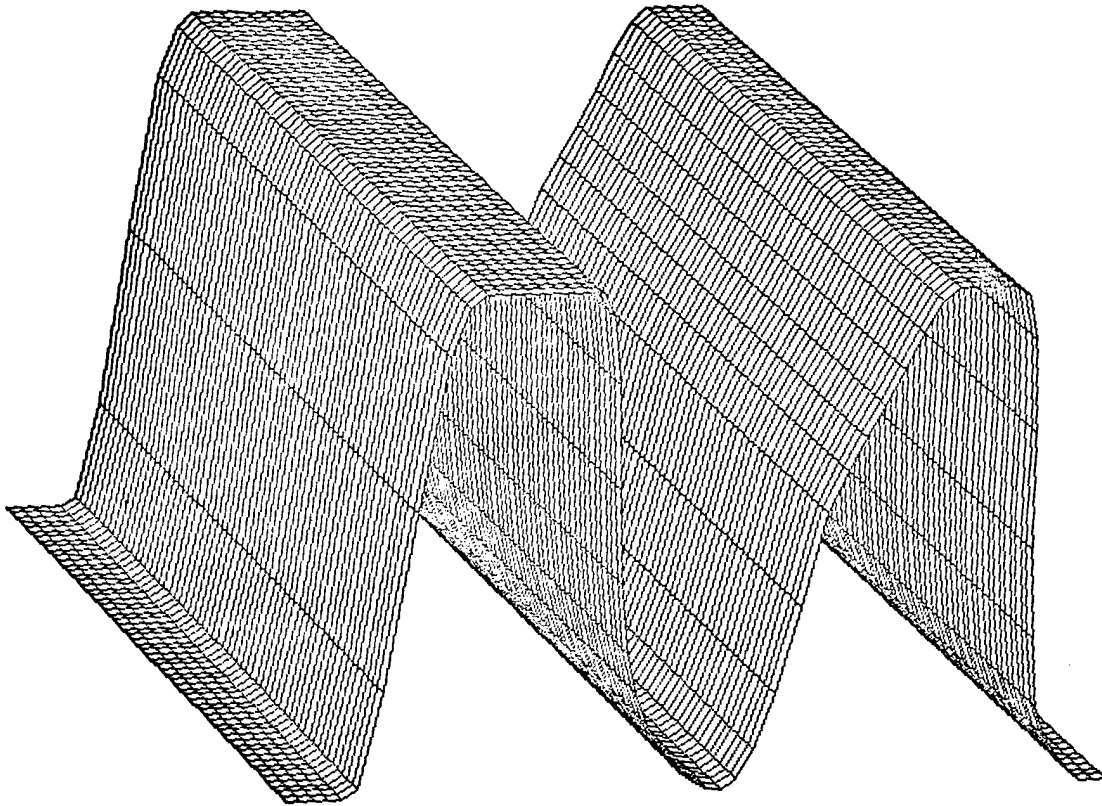


Fig. 27. Lamb wave transmission through the bond of the alumina flexure-strength specimen illustrated in Fig. 25, showing detection of the simulated nonbond.

Figure 28 shows the results obtained with Lamb waves on that second specimen. There is a small region near the center of the braze joint where the Lamb wave transmission dips noticeably. We then switched to a pulse-echo mode; that is, we monitored the amplitude of the Lamb waves reflected by the interface using first one and then the other transducer as a transmitter. In each case, an enhanced reflection was found at the location of the dip in Fig. 28. This eliminates problems in the ceramic (e.g., cracking) away from the interface as the source of the anomaly and establishes that some variation in the region of the interface engendered the signal.

Radiographic examination of the specimen revealed no detectable flaws in the interface.

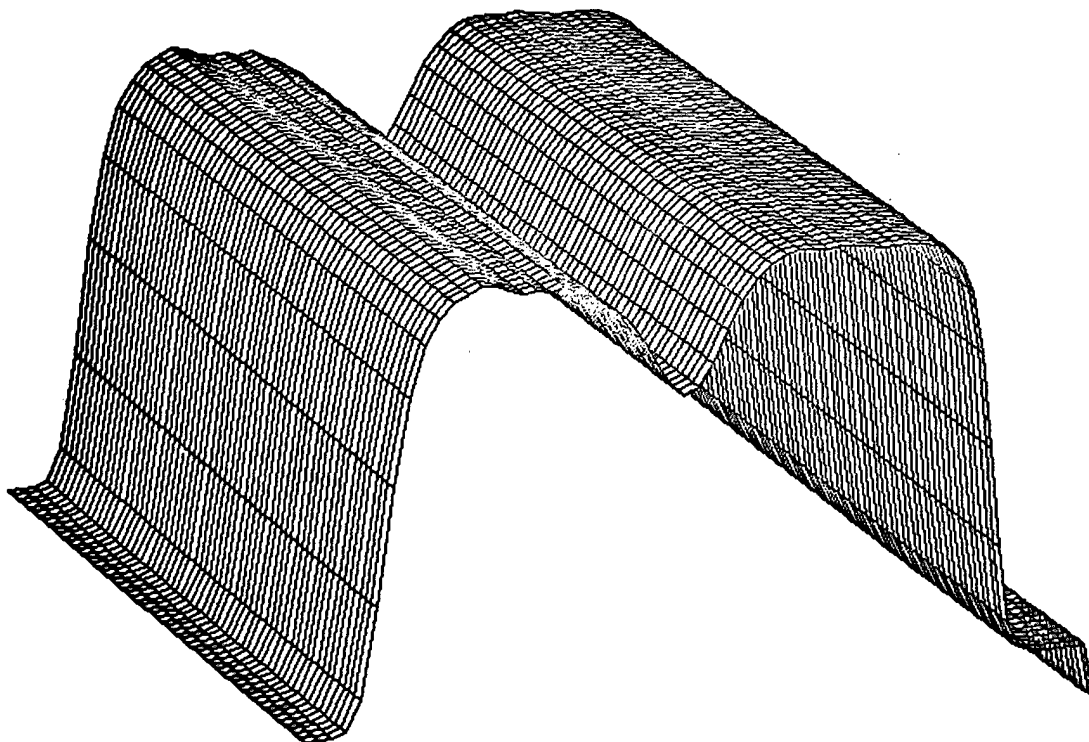


Fig. 28. Lamb wave transmission through the bond of another PSZ flexure-strength specimen having a simulated nonbond, showing a small indication near the center of the bond.

These results indicate that Lamb waves are sensitive to some characteristics of the braze joint that conventional nondestructive test techniques do not detect. The specimen will be cut into flexure bars and subjected to bend tests to try to develop a correlation between the inspection results and the strength of the bond.

#### SUMMARY AND CONCLUSIONS

The small critical flaw size in conventional monolithic structural ceramics has placed severe burdens on the development of ultrasonic NDE techniques for ensuring the integrity of parts fabricated of these materials. Since the detectability of a given flaw increases as the wavelength of the interrogating radiation decreases (at least until the wavelength becomes comparable to the flaw size), frequencies of at least 50 MHz are required to detect critical flaws in many ceramics. In addition, unlike structural metals, the interrogating ultrasonic radiation

must be focused in order to ensure that a sufficiently large fraction of the incident radiation is intercepted by the flaw. The necessity of introducing focused radiation through the surface of the ceramic part (mandated by the desire to have rapid scanning capability) leads to severe spherical aberration of the ultrasonic beam within the part and limits the depth to which critical flaws can be detected to a few millimeters. Within these constraints, however, we have had considerable success in detecting flaws in the critical size range.

Since a common flaw type in ceramics is a quasi-spherical void or inclusion, a model was discussed for the scattering of ultrasonic waves from spheres. This model was applied to scattering from natural flaws in partially stabilized zirconia, and a flaw diameter of about 25  $\mu\text{m}$  was inferred.

Since the attenuation behavior of structural ceramics for ultrasonic waves is critical to determining the minimum detectable flaw size, a technique was developed that permits a material transfer curve, or attenuation versus frequency response, to be determined for any ceramic. This process corrects for all known material-independent losses, such as diffraction (beam spread), acoustic impedance mismatches at the surfaces of the part and at internal interfaces, and frequency-dependent coupling losses between the transducer and the surface of the part. The transfer curve is highly sensitive to changes in the microstructure of the ceramic, and examples of the differences in TZP and PSZ ceramics were given.

Since current heat engine designs contain components consisting of a ceramic material clad on a metallic substrate to achieve both high-temperature tolerance and toughness, additional evaluation techniques were required for assurance of the quality of a ceramic-metal bond. The braze layer in such a bond is typically of the order of 60  $\mu\text{m}$  thick, and, if nonbonding occurs, the tester needs to determine which face of this layer is nonbonded. For high-frequency ceramics, this condition could be determined by interrogating the bond with frequencies sufficiently high ( $\approx 100$  MHz) to resolve the layer thickness. Examples were given of the detection of 100- $\mu\text{m}$ -diam pores in the braze layer using this approach. When the ceramic host would not support ultrasonic waves of such a high

frequency, however, advanced signal processing techniques had to be developed to recover the desired information. An example was given of the identification of the correct interface in a braze joint in PSZ.

Although no specimens were available for study that contained manufactured flaws of known size in the bond layer of a ceramic-to-metal joint, a sample containing such flaws at the interface between a P-leg and an N-leg of a silicon-germanium specimen was obtained. The detection of a 125- $\mu\text{m}$  flaw was demonstrated, and the signal-to-noise ratio for this flaw indicates that considerably smaller flaws could be detected reliably.

In testing joints between ceramic coupons, it is often not possible to introduce the ultrasound at normal incidence to the bond because of specimen geometry. In such cases, angle-beam testing techniques can be used, but it is difficult to maintain transducer focus across the full width of the interface. In addition, angle-beam testing was found to be unduly sensitive to slight vertical misalignment of the two plates being joined. For these configurations, a second approach, one using plate or Lamb waves, was studied. Two transducers were used: one launches a Lamb wave in either plate of the specimen, and the second receives the wave in the other plate after propagation through the joint. The amplitude of the received wave is an indicator of the condition of the bond in the region between the transducers.

Dispersion curves for Lamb waves in alumina coupons were calculated, and a system was assembled for evaluation of the bond by these waves. The system was shown to detect a simulated nonbond easily, and inspection of a second sample by this technique revealed an indication in the bond area that was not detected by angle-beam testing.

While bonds between ceramic components and between ceramic and metallic parts present a number of difficult problems to the researcher, considerable progress has been made in developing techniques that will allow the integrity of these bonds to be assured nondestructively. Many problems remain, however. Perhaps the most common criticism directed at ceramic inspection is that it is slow in comparison with inspection of metals. Rarely is it noted, however, that most ceramic inspection systems do not approach the state of the art in data acquisition rates. This is understandable, since technique development is of more importance at the

present time in ceramic evaluation than it is for metals, where the techniques have a long history of development. Inspection rates will undoubtedly increase for ceramic evaluation as the approaches become established.

A more fundamental limitation in ceramic or ceramic-joint evaluation is related to the requirement for focused radiation. The severe aberrations introduced into the beam by propagation through the sample surface limit the depth to which effective focus can be maintained. This depth can be increased somewhat by increasing transducer frequency and focal length, but the relationship does not appear to favor this approach. Some form of this limitation will likely persist in the foreseeable future.

The techniques presented here do not exhaust the potential tools for ceramic-joint evaluation by a wide margin. For example, some form of direct interface wave (i.e., a wave that propagates in the bond material but not in the ceramic) may possibly yield information about the strength of the bond itself. This result would indeed be valuable, since there is currently no known correlation between bond strength and measurable acoustic properties.

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